

1) FT systems

$$F(x) \rightarrow \boxed{\text{FT}} \rightarrow \mathcal{F}(F(x))$$

a) Linear?

FF linear: The output can be written as a summation (integral) of individual responses to unit impulses.

$$\text{--or-- } af(x) + bg(x) \rightarrow \boxed{\text{FT}} \rightarrow a\mathcal{F}(f(x)) + b\mathcal{F}(g(x))$$

True by properties of Fourier Transform.

b)

Can you specify a transfer function for this system?

That is: Is the FT both linear & shift invariant?

If the system is shift invariant then

$$\text{FT} \{ f(x-x_0, y-y_0) \} = F(u-x_0, v-y_0)$$

But from the shift theorem we know

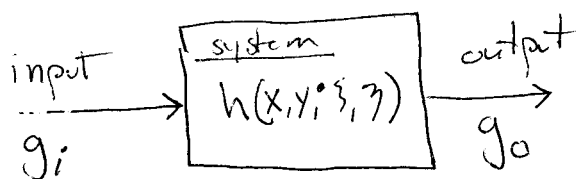
$$\text{FT} \{ f(x-x_0, y-y_0) \} = F(u, v) e^{-i2\pi(ux_0 + vy_0)}$$

$$\therefore F(u-x_0, v-y_0) = F(u, v) e^{-i2\pi(ux_0 + vy_0)}$$

So this is a contradiction, because we know a valid Transform exists (eg $\text{sinc}(au)$) that violates this.

$\therefore$ FT is not shift invariant $\Rightarrow$ ~~not~~ transfer function.

#2



$h(x, y; \xi, \eta) \equiv$ The response of the system at point x, y in the output space to a δ function input at (ξ, η) of the input space

The impulse response

If $h(x, y; \xi, \eta)$ is linear the output can be written as a summation (integral) of individual responses to unit impulses

$$g_o(x, y) = \iint_{-\infty}^{\infty} g_i(\xi, \eta) h(x, y; \xi, \eta) d\xi d\eta$$

// If the system is shift invariant, the impulse response $h(x, y; \xi, \eta)$ depends only on the distances $(x - \xi)$ and $(y - \eta)$. And we take a simpler notation

$$h(x, y; \xi, \eta) = h(x - \xi, y - \eta)$$

Now the output can be written

$$g_o(x, y) = \iint_{-\infty}^{\infty} g_i(\xi, \eta) h(x - \xi, y - \eta) d\xi d\eta$$

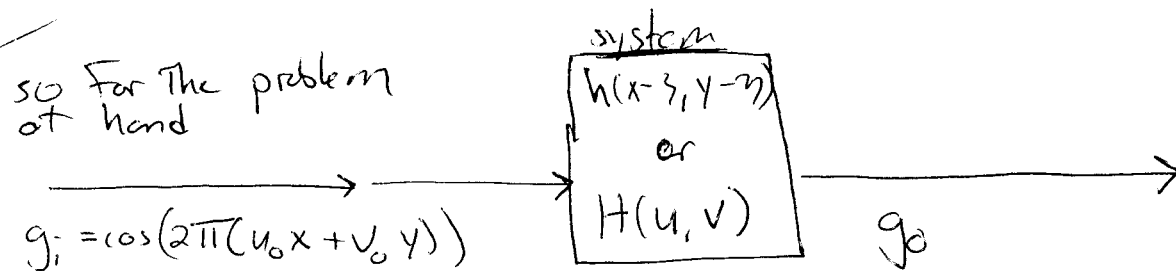
which is the two dimensional convolution, written in a simpler notation

$$g_o(x, y) = g_i(x, y) \otimes h(x, y)$$

The reason we like linear, shift-invariant systems is now revealed when we take the Fourier Transform of the input/output relation:

$$\mathcal{FT}(g_o(x, y) = g_i(x, y) \otimes h(x, y))$$

$$\Rightarrow G_o(u, v) = G_i(u, v) H(u, v)$$



$$G_o(u, v) = G_i(u, v) H(u, v)$$

$$\begin{aligned} G_i(u, v) &= \iint \cos(2\pi(u_0 x + v_0 y)) e^{-i2\pi(ux + vy)} dx dy \\ &= \iint \frac{1}{2} \left[e^{i2\pi(u_0 x + v_0 y)} + e^{-i2\pi(u_0 x + v_0 y)} \right] e^{-i2\pi(ux + vy)} dx dy \\ &= \iint \frac{1}{2} e^{i(2\pi u_0 x + 2\pi v_0 y - 2\pi u x - 2\pi v y)} dx dy + \frac{1}{2} \iint e^{-i(2\pi u_0 x + 2\pi v_0 y + 2\pi u x + 2\pi v y)} dx dy \\ &= \iint \frac{1}{2} e^{-i2\pi[(u-u_0)x + (v-v_0)y]} dx dy + \iint \frac{1}{2} e^{-i2\pi[(u+u_0)x + (v+v_0)y]} dx dy \\ &= \frac{1}{2} \delta(u-u_0, v-v_0) + \frac{1}{2} \delta(u+u_0, v+v_0) \end{aligned}$$

$$G_0(u, v) = G_1(u, v) \times H(u, v)$$

⇒ Clearly the output is the same frequency as input. with a different Amplitude & phase

$$G_0(u,v) = \frac{1}{2} H(u_0, v_0) \delta(u-u_0, v-v_0) + \frac{1}{2} H(-u_0, -v_0) \delta(u+u_0, v+v_0)$$

$$g_0(x, y) = f^{-1}(G_0(u, v))$$

$$= \frac{1}{2} H(u_0, v_0) e^{i2\pi(u_0 x + v_0 y)} + \frac{1}{2} H(-u_0, -v_0) e^{-i2\pi(u_0 x + v_0 y)}$$

$$g(x, y) = |M| \cos[2\pi(u_0 x + v_0 y) - \phi]$$

$$\omega/ \quad |M| = 2 \sqrt{\frac{1}{2} H(u_0, v_0) H(-u_0, -v_0)}$$

$$\phi = \tan^{-1} \left(\frac{i \left[\frac{1}{2} H(-u_0, v_0) + \frac{1}{2} H(+u_0, v_0) \right]}{\left[\frac{1}{2} H(-u_0, -v_0) - \frac{1}{2} H(u_0, v_0) \right]} \right)$$

See details on next page..

Details

$$\Rightarrow Ae^{i\theta x} + Be^{-i\theta x} = D \sin(\theta x) + E \cos(\theta x)$$
$$= D \left(\frac{e^{i\theta x} - e^{-i\theta x}}{2i} \right) + E \left(\frac{e^{i\theta x} + e^{-i\theta x}}{2} \right)$$
$$= \left(\frac{D}{2i} + \frac{E}{2} \right) e^{i\theta x} + \left(\frac{-D}{2i} + \frac{E}{2} \right) e^{-i\theta x}$$

$$\therefore D + iE = 2iA \Rightarrow -2iA + iE = -D$$

$$-D + iE = 2iB \Rightarrow -2iA + iE + iE = 2iB$$

$$-2iA + 2iE = 2iB$$

$$\boxed{E = B + A}$$

$$D = 2iA - iE = 2iA - i(B + A)$$

$$= 2iA - iB - iA$$

$$= iA - iB$$

$$\boxed{D = i(A - B)}$$

now

$$D \sin(\theta x) + E \cos(\theta x) = |M| \cos(\theta x - \phi)$$

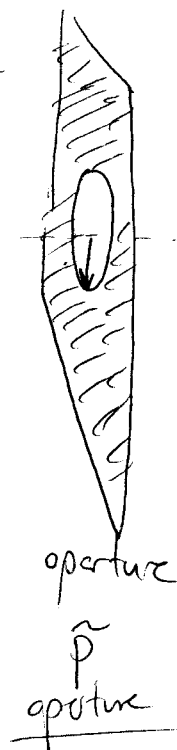
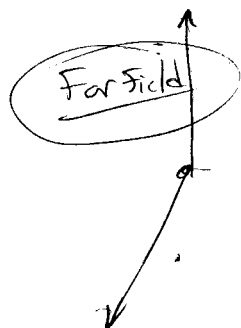
$$|M| = \sqrt{D^2 + E^2} = \sqrt{i^2(A-B)^2 + (B+A)^2} = 2\sqrt{AB}$$

$$\phi = \tan^{-1} \left(\frac{E}{D} \right) = \tan^{-1} \left(\frac{B+A}{-i(B-A)} \right) = \tan^{-1} \left(\frac{i(B+A)}{B-A} \right)$$

#3 a) $T(u, v)$ for 

First lets find $T(u, v)$ for 

$$\tau(x, y) = \begin{cases} 1 & r < r_0 \\ 0 & \text{otherwise} \end{cases}$$



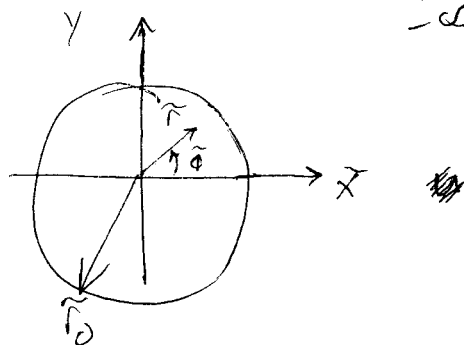
Far Field approximation

where distance is assumed to be linear in $\tilde{x}$ & $\tilde{y}$

(i.e. $\sqrt{1+x} \approx 1 + \frac{1}{2}x$)

we find that the observed field at P' due to a point source at P w/ an aperture τ is given by

$$E(u, v) \propto T(u, v) = \iint_{-\infty}^{\infty} \tau(\tilde{x}, \tilde{y}) e^{-i2\pi(u\tilde{x} + v\tilde{y})} d\tilde{x} d\tilde{y}$$



$$\tilde{x} = \tilde{r} \cos \tilde{\phi}$$

$$\tilde{y} = \tilde{r} \sin \tilde{\phi}$$

$$T(u, v) = \iint_{-\infty}^{\infty} \begin{pmatrix} 1 & \text{inside circle} \\ 0 & \text{outside circle} \end{pmatrix} e^{-i2\pi(u\tilde{r}\cos\tilde{\phi} + v\tilde{r}\sin\tilde{\phi})} \tilde{r} d\tilde{r} d\tilde{\phi}$$

Further
exploit circular symmetry

$$\text{let } \rho \equiv \sqrt{u^2 + v^2}$$

$$u = \rho \cos \phi'$$

$$v = \rho \sin \phi'$$

$$T(\rho, \phi') = \int_0^{\rho_0} \int_0^{2\pi} 1 \cdot e^{-i2\pi(\rho\tilde{r}\cos\tilde{\phi}\cos\phi' + \rho\tilde{r}\sin\tilde{\phi}\sin\phi')} \tilde{r} d\tilde{r} d\tilde{\phi}$$

$$= \int_0^{\rho_0} \int_0^{2\pi} e^{-i2\pi(\rho\tilde{r})(\cos\tilde{\phi} - \phi')} \tilde{r} d\tilde{r} d\phi$$

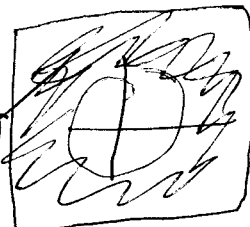
Note: (eg) 6.45 $J_0(w) = \frac{1}{2\pi} \int_0^{2\pi} \exp[-i w \cos(\tilde{\phi} - \phi')] d\tilde{\phi}$

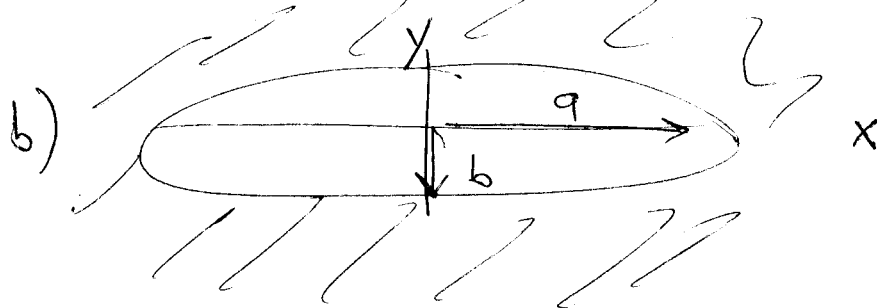
$$\int_0^x \xi J_0(\xi) d\xi = x J_1(x)$$

$$\pi(\rho, \phi') = 2\pi \int_0^{\rho_0} \tilde{r} J_0(2\pi\rho\tilde{r}) d\tilde{r} = \frac{\rho_0}{\rho} J_1(2\pi\rho\rho_0)$$

$$T(u, v) = \frac{\rho_0}{\sqrt{u^2 + v^2}} J_1(2\pi\sqrt{u^2 + v^2} \rho_0) \quad \text{This is for}$$

$$T(u, v) \text{ of } \left[\text{circle} \right] = \sqrt{u^2 + v^2} - \frac{\rho_0}{\sqrt{u^2 + v^2}} J_1(2\pi\sqrt{u^2 + v^2} \rho_0)$$





$$f(x, y) = \begin{cases} 1 & \text{inside ellipse} \\ 0 & \text{out} \end{cases} \Rightarrow F(u, v)?$$

equation for ellipse $r^2 = \frac{x^2}{a^2} + \frac{y^2}{b^2}$

$$x = a \cos \theta \quad y = b \sin \theta$$

$$F(u, v) = \iint_{\text{ellipse}} e^{-i2\pi(ux+vy)} dx dy$$

change of variables

Note in general: (see attached copies from Hildebrand)

$$\iint_D w(x, y) dx dy = \iint_{D^*} w(u, v) \underbrace{\left| \frac{\partial(x, y)}{\partial(u, v)} \right|}_{\text{Jacobian}} du dv$$

$$\left| \frac{\partial(x, y)}{\partial(r, \theta)} \right| = \begin{vmatrix} \frac{\partial x}{\partial r} & \frac{\partial x}{\partial \theta} \\ \frac{\partial y}{\partial r} & \frac{\partial y}{\partial \theta} \end{vmatrix} = \begin{vmatrix} a \cos \theta & -a \sin \theta \\ b \sin \theta & b \cos \theta \end{vmatrix} = ab$$

so

$$F(u, v) = \int_0^{2\pi} d\theta \int_0^1 ab r dr e^{-i2\pi(ua \cos \theta + vb \sin \theta)}$$

$$\begin{aligned} \text{let } u a &= \rho \cos \phi \\ v b &= \rho \sin \phi \end{aligned}$$

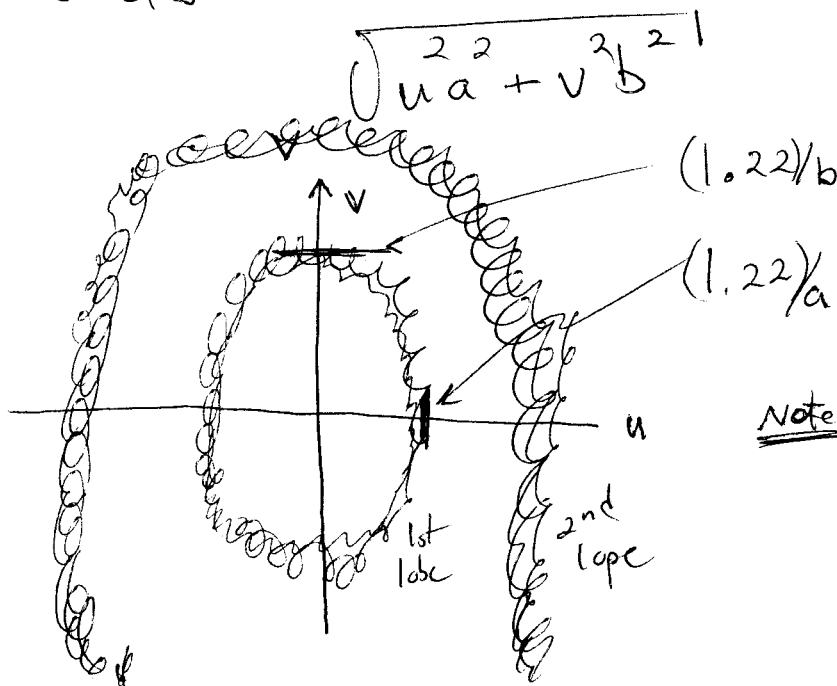
$$\rho^2 = u^2 a^2 + v^2 b^2$$

$$\begin{aligned} F(\rho, \phi) &= \int_0^{2\pi} \int_0^1 a b r e^{-i 2\pi (r \rho \cos \theta \cos \phi + r \rho \sin \theta \sin \phi)} dr d\phi \\ &= a b \int_0^{2\pi} \int_0^1 r e^{-i 2\pi \rho r \cos(\theta - \phi)} dr d\phi \end{aligned}$$

$$= 2\pi a b \int_0^1 r J_0(2\pi \rho r) dr$$

$$= a b \frac{J_1(2\pi \rho)}{\rho}$$

$$F(u, v) = a b \frac{J_1(2\pi \sqrt{u^2 a^2 + v^2 b^2})}{\sqrt{u^2 a^2 + v^2 b^2}}$$



Note: the First lobe of FT is also an ellipse oriented 90° wrt the original ellipse

wavelength	n
0.40	1.47011
0.42	1.46809
0.44	1.46634
0.46	1.46483
0.48	1.46350
0.50	1.46232
0.52	1.46128
0.54	1.46034
0.56	1.45949
0.58	1.45873
0.60	1.45803
0.62	1.45739
0.64	1.45681
0.66	1.45626
0.68	1.45576
0.70	1.45529

relection coeff.
-0.19032
-0.18966
-0.18908
-0.18859
-0.18815
-0.18776
-0.18741
-0.18710
-0.18682
-0.18657
-0.18634
-0.18613
-0.18594
-0.18575
-0.18559
-0.18543

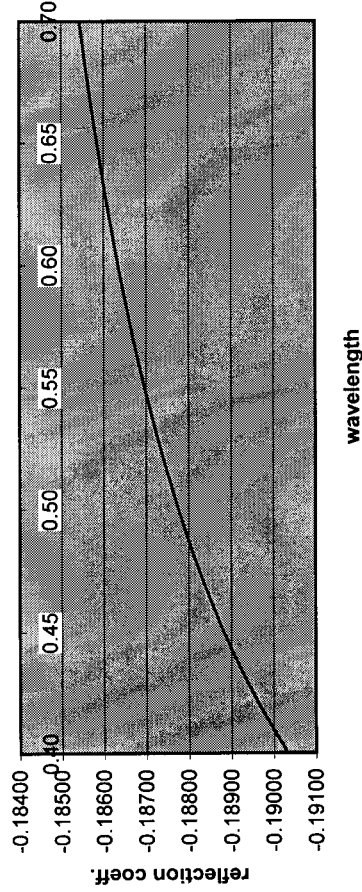
average: -0.18729

0.0351

Energy Refl. Coeff = $|p|^2$

Refl. Coeff
0.0362
0.0360
0.0358
0.0356
0.0354
0.0353
0.0351
0.0350
0.0349
0.0348
0.0347
0.0346
0.0346
0.0345
0.0344
0.0344

relection coeff. p_r



flux density $\equiv$ flux per unit area at a point on a surface $[W/m^2]$

3a) Field Reflection

$$p_r = \frac{1 - n(\lambda)}{1 + n(\lambda)}$$

or $R = p_r^2$

Energy Reflection

both accepted

3b) a couple of things were accepted.

$\langle p_r \rangle \equiv \text{avege of } p_r \approx -0.187$

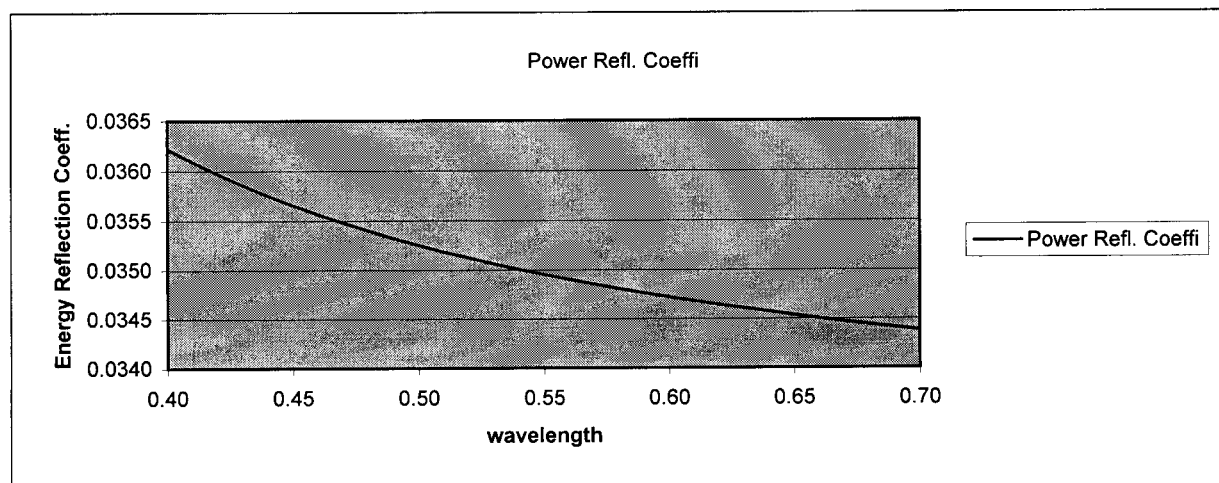
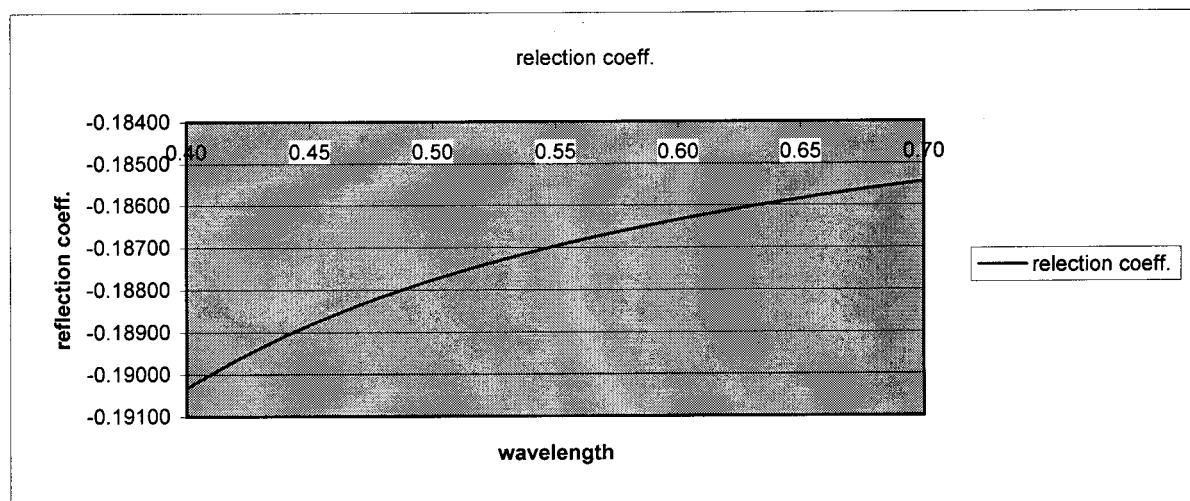
or

$\langle R \rangle \equiv \text{avege power reflected} \approx 3.5\%$

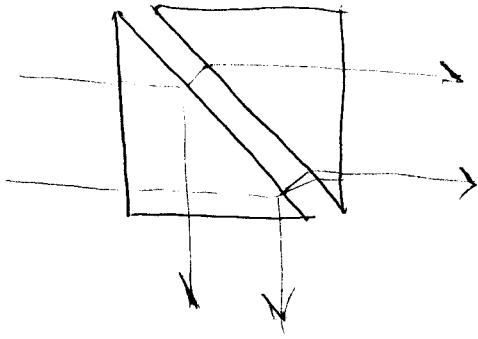
i.e. what would actually be observed w/ your eye or a sensor

This is actually what was wanted in the question.

wavelength	n	relection coeff.	Power Refl. Coeffi
0.40	1.47011	-0.19032	0.0362
0.42	1.46809	-0.18966	0.0360
0.44	1.46634	-0.18908	0.0358
0.46	1.46483	-0.18859	0.0356
0.48	1.46350	-0.18815	0.0354
0.50	1.46232	-0.18776	0.0353
0.52	1.46128	-0.18741	0.0351
0.54	1.46034	-0.18710	0.0350
0.56	1.45949	-0.18682	0.0349
0.58	1.45873	-0.18657	0.0348
0.60	1.45803	-0.18634	0.0347
0.62	1.45739	-0.18613	0.0346
0.64	1.45681	-0.18594	0.0346
0.66	1.45626	-0.18575	0.0345
0.68	1.45576	-0.18559	0.0344
0.70	1.45529	-0.18543	0.0344
average:		-0.18729	0.0351



#6



$$n = 1.6$$

$$\lambda = 632.8 \text{ nm} \quad (\pi \text{ polarized})$$

find size air gap that produces a Transmission coefficient of 50%

energy Transmission

$$T_{\pi} = 2e^{-2d/\delta} (1 - \cos 2\phi_{\pi}) \quad \text{eq 2.84}$$

$$d \equiv \text{gap}$$

$$\delta \equiv \text{skin depth}$$

$$\phi_{\pi} - \pi \text{ phase angle}$$

$$d = -\frac{\delta}{2} \ln \left[\frac{T_{\pi}}{2(1 - \cos 2\phi_{\pi})} \right]$$

$$\delta \text{ skin depth from eq 2.82}$$

$$\delta = \frac{\lambda_0}{2\pi \sqrt{n^2 \sin^2 \theta - n'^2}}$$

$$\tan \frac{\phi_{\pi}}{2} = \frac{n^2}{n'^2} \gamma \Rightarrow \phi_{\pi} = 2 \tan^{-1} \left(\frac{n^2}{n'^2} \gamma \right) \quad (\text{follows from 2.83})$$

γ - describes evanescent wave
(analog to wave vector that decays)

$$\gamma = \frac{\sqrt{n^2 \sin^2 \theta - n'^2}}{n \cos \theta}$$

so all together

$$d = -\frac{\delta}{2} \ln \left[\frac{T_{\pi}}{2(1 - \cos 2\phi_{\pi})} \right]$$

$$\delta = \frac{\lambda_0}{2\pi(n^2 \sin^2 \theta - n'^2)^{1/2}}$$

$$\phi_{\pi} = 2 \tan^{-1} \left(\frac{n^2}{n'^2} \gamma \right)$$

$$\gamma = \frac{\sqrt{n^2 \sin^2 \theta - n'^2}}{n \cos \theta}$$

with

$$\theta = \frac{\pi}{4} \quad T_{\pi} = .50 \quad n' = 1 \quad n = 1.6$$

$$\lambda = 638.8 \text{ nm}$$

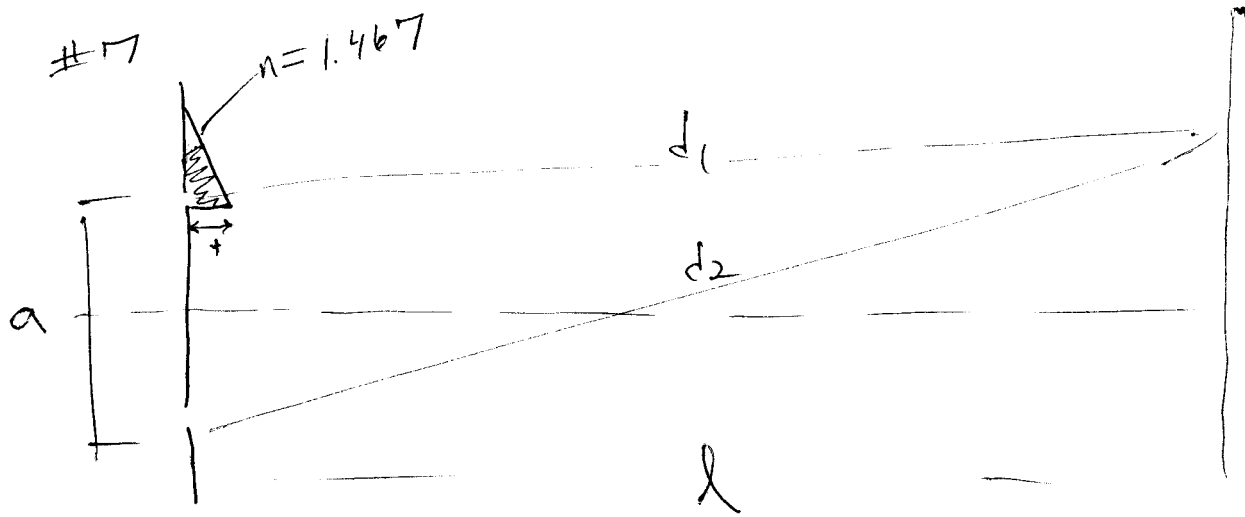
//

$$\gamma = .4677$$

$$\delta = 190.3301 \text{ nm}$$

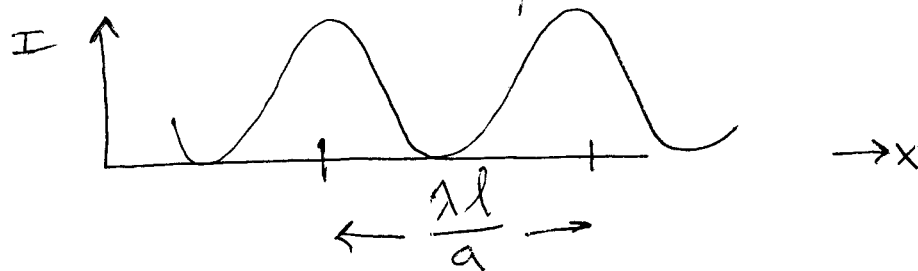
$$\phi_{\pi} = 1.7499 \text{ rad}$$

$$\boxed{d = 194.82 \text{ nm}}$$



- central max shifts by 3.4 Fringes to one side
- find max Thickness of wedge

we know that for 2 point interference



so we have a distance shift of $3.4 \left(\frac{\lambda l}{a} \right)$

at maximum the phase difference between the paths is zero.

phase difference:

$$\phi = \frac{2\pi}{\lambda} (\text{OPL}_1 - \text{OPL}_2)$$

$$= \frac{2\pi}{\lambda} [(nt + d_1 - t) - d_2]$$

$$= \frac{2\pi}{\lambda} [(n-1)t + d_1 - d_2]$$

$$\begin{aligned}
 d_1 &= \ell^2 + \left(\frac{a}{2} - x\right)^2 \\
 d_1 &= \ell \sqrt{1 + \left(\frac{a/2 - x}{\ell}\right)^2} \approx \ell \left[1 + \frac{1}{2} \left(\frac{a/2 - x}{\ell}\right)^2\right] \\
 d_2 & \approx \ell \left[1 + \frac{1}{2} \left(\frac{a/2 + x}{\ell}\right)^2\right]
 \end{aligned}$$

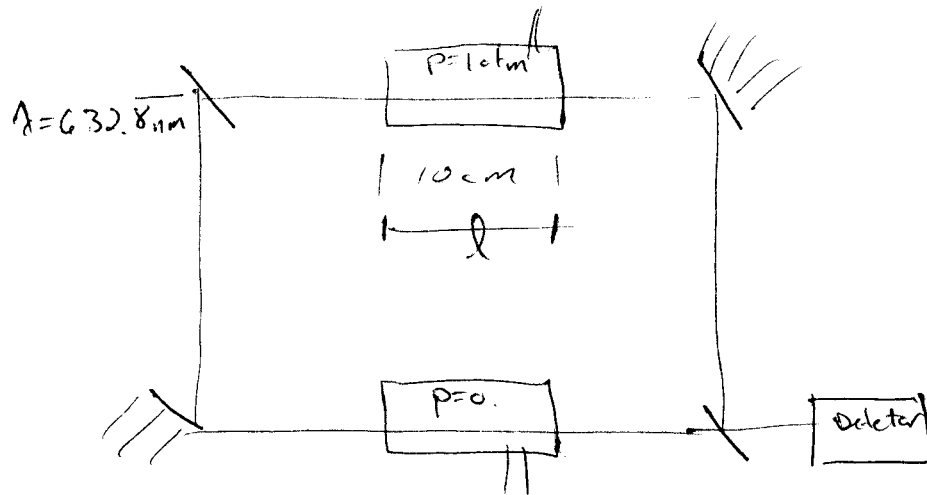
$$\phi = \frac{2\pi}{\lambda} \left[(n-1)\ell + \overset{\substack{\uparrow \\ d_1 - d_2}}{\frac{ax}{\ell}} \right]$$

so
at $x = 3.4 \frac{\lambda \ell}{a}$; $\phi = 0$

$$\phi = \frac{2\pi}{\lambda} \left[(n-1)\ell + \frac{a}{\ell} (3.4) \frac{\lambda \ell}{a} \right] = 0$$

$$\boxed{\cancel{\text{path difference}} + \frac{3.4 \lambda}{(n-1)} = \frac{3.4 \lambda}{1.467} = 7.2805 \lambda}$$

8 KF 5.37



100 minima are observed as cell is evacuated. ← gave Through 2π 100 times
 n For gas is $n = 1 + A p$ - pressure

$$\phi = \frac{2\pi}{\lambda_0} (\text{OPD})$$

← optical path difference

$$= \frac{2\pi}{\lambda_0} (n-1) l$$

$$\phi = \frac{2\pi}{\lambda_0} [(1 + A p) - 1] l$$

$$\frac{\phi \lambda_0}{2\pi l p} = A$$

$$A = \frac{2\pi \cdot 100 \cdot \lambda_0}{2\pi l p} = 6.328 \times 10^{-4} \frac{1}{\text{atm}}$$

if strict:

$$A = \frac{199\pi \cdot \lambda_0}{2\pi l p} = 6.296 \times 10^{-4} \frac{1}{\text{atm}}$$

to be strict:



The first minimum occurs after π shift, then every 2π after that

so actually ϕ is

$$200\pi - \pi = 199\pi$$

accepted answer